

### A. Zero temperature

Susceptibility can be easily calculated at  $T = 0$  for any momentum and frequency. Here for simplicity we drop the frequency.

$$\chi_0(q) = \frac{1}{2\pi} \sum_k \int d\omega \frac{1}{(i\omega - \varepsilon_{k+q})(i\omega - \varepsilon_k)} \quad (5)$$

$$= -2 \sum_k \frac{\Theta(-\varepsilon_k)}{\varepsilon_k - \varepsilon_{k+q}} \quad (6)$$

$$= \frac{1}{2\pi^2} \int \frac{\Theta(-\varepsilon_k) k dk d\phi}{kq \cos \phi + q^2/2} \quad (7)$$

$$= \frac{1}{\pi^2} \int \frac{\Theta(-\varepsilon_k) k dk d\phi}{kq(z + z^{-1}) + q^2} \quad (8)$$

$$= \frac{1}{\pi^2} \int_0^{k_F} \frac{dk}{q} \int_C \frac{dz/i}{z^2 + zq/k + 1}. \quad (9)$$

Here  $C$  is the unit circle and  $z = e^{i\phi}$ . The poles of the integrand are  $z_{\pm} = -q/(2k) \pm \sqrt{q^2/(2k)^2 - 1}$ . For  $k < q/2$  the  $z_+$  pole is inside unit circle and  $z_-$  outside. Therefore,

$$\chi_0(q) = \frac{1}{\pi} \int_0^{\min(k_F, q/2)} \frac{dk}{q} \frac{1}{\sqrt{q^2/(2k)^2 - 1}} \quad (10)$$

$$= \frac{1}{\pi q} \sqrt{q^2/4 - k^2} \Big|_{\min(k_F, q/2)}^0, \quad (11)$$

and therefore,

$$\chi_0(q < 2k_F) = \frac{1}{2\pi} \quad (12)$$

$$\chi_0(q > 2k_F) = \frac{1}{2\pi} \left( 1 - \sqrt{1 - 4k_F^2/q^2} \right) \quad (13)$$

We see that zero-temperature susceptibility is completely flat for  $q < 2k_F$ . This remains true (in a reduced interval of momenta) also at finite frequency.

### B. Finite temperature

Let us now determine the form of susceptibility at finite temperature, first assuming that  $(k_F q, \Omega) \ll T$ . For brevity assume that  $\Omega = 0$ .

Then from Eq. (??) after performing summation over  $\omega_n$ ,

$$\chi(q) = - \sum_k \frac{n_F(\varepsilon_k) - n_F(\varepsilon_{k+q})}{\varepsilon_k - \varepsilon_{k+q}} \quad (14)$$

$$= - \sum_k n'_F(\varepsilon_k) + \frac{1}{2} n''_F(\varepsilon_k) (\varepsilon_{k+q} - \varepsilon_k) + \frac{1}{6} n'''_F(\varepsilon_k) (\varepsilon_{k+q} - \varepsilon_k)^2 + \dots \quad (15)$$

$$= - \sum_k n'_F(\varepsilon_k) + \frac{1}{2} n''_F(\varepsilon_k) (\mathbf{kq} + q^2/2) + \frac{1}{6} n'''_F(\varepsilon_k) (\mathbf{kq} + q^2/2)^2 + \dots \quad (16)$$

$$= - \int \frac{k dk d\phi}{(2\pi)^2} \left[ n'_F(\varepsilon_k) + \frac{1}{2} n''_F(\varepsilon_k) (\mathbf{kq} + q^2/2) + \frac{1}{6} n'''_F(\varepsilon_k) (\mathbf{kq} + q^2/2)^2 + \dots \right] \quad (17)$$

$$= - \int \frac{d\varepsilon_k}{2\pi} \left[ n'_F(\varepsilon_k) + \frac{1}{2} n''_F(\varepsilon_k) q^2/2 + \frac{1}{6} n'''_F(\varepsilon_k) ((\varepsilon_k + \mu)q^2 + q^4/4) + \dots \right] \quad (18)$$

$$= \frac{1}{2\pi} \quad (19)$$